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On the Existence of a Riemannian Manifolds at a Given Connection

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Abstract

We give necessary and sufficient conditions for a linear connection without torsion to come from a Riemannian metric. The nullity space and the image space of the curvature are involved in the formulation of the results.

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Introduction

A Riemannian manifold defines a unique linear connection without torsion and conservative. By studying the properties of the space of nullity and images of the curvatures, we reduce the resolution of the inverse problem to a system of linear equations with some properties of the curvatures quite simple to verify.

The formalism used is that of [1] and [2] and the formulas are those of [3], the method following [4].

Canonical Connection of a Riemannian Manifold

Let M be a n dimensional differentiable paracompact manifold and of class \mathcal{C}^{∞} , J is the natural tangent structure of the tangent bundle: TM \rightarrow M. The derivation d_{J} [5] is defined by d_{J} =[i_{J} , d].

Definition 1

We call the Riemannian manifold, [1], the couple (M, E):

M a differentiable manifold

• E a function of $TM = TM - \{0\}$ in \mathbb{R}^+ , with E(0)=0, C^{∞} on TM, C^2 on the null section, and homogeneous of degree two such that dd_JE has a maximum rank

The application E is called energy function, the fundamental scalar 2-form Ω =dd,E defines a spray S [6]:

$$i_{s}dd_{J}E = -dE, \tag{2.1}$$

the derivation i_s being the inner product with respect to S. The vector 1-form Γ =[J, S] is called canonical connection. The connection Γ thus defined is an almost product structure: Γ^2 =I, I is the vector 1-form identity. By asking

$$h = \frac{1}{2}(I + \Gamma)$$
 and $v = \frac{1}{2}(I - \Gamma)$,

h is the horizontal projector, projector of the proper subspace at the

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eigenvalue +1, v the vertical projector corresponding to the eigenvalue-1. The curvature of Γ is defined by

$$R = \frac{1}{2}[h,h] \tag{2.2}$$

The vector 2-form R is also called the Nijenhuis tensor of h. The scalar 2-form allows to define a metric g on the vertical bundle by

$$g(JX, JY) = \Omega (JX, Y)$$

for all $X, Y \in \chi(TM)$, where $\chi(TM)$ denotes the set of vector fields on TM.

There exists, [2], one and only one metric lift ${\sf D}$ of the canonical connection such that:

- 1. DJ=0;
- 2. DC=v, (C being the Liouville canonical field on TM);
- 3. DΓ=0;
- 4. Dg=0.

The D connection is called the Cartan connection. We have

$$D_{JX}JY = [J, JY]X, \quad D_{hX}JY = [h, JY]X.$$
 (2.3)

With the linear connection D, we associate a curvature

 $\Re(X,Y)Z = D_{hX}D_{hY}JZ - D_{hY}D_{hX}JZ - D_{[hX,hY]}JZ,$

for all $X, Y, Z \in \chi(TM)$. The relation between the curvatures \Re and R is [2]:

$$\Re(X,Y)Z = J[Z,R(X,Y)] - [JZ,R(X,Y)] + R([JZ,X],Y) + R(X,[JZ,Y]).$$
(2.4)

In particular,

$$\Re(X,Y)S = -R(X,Y). \tag{2.5}$$

In natural local coordinates on an open set U of M, $(x^i, y^i) \in TU$, i, $j \in \{1,...,n\}$. The energy function is written [1] p.330

$$E = \frac{1}{2}g_{ij}(x^{1},...,x^{n})y^{i}y^{j}$$

where $g_{ij}(x^1,...,x^n)$ are symmetric positive functions such that the matrix $(g_{ii}(x^1,...,x^n)_i)$ is invertible. And the relation $i_S dd_J E = -dE$ gives the spray S

$$S = y^{i} \frac{\partial}{\partial x^{i}} - 2G^{i}(x^{1}, ..., x^{n}, y^{1}, ..., y^{n}) \frac{\partial}{\partial y^{i}}$$

with

$$G_k = \frac{1}{2} y^i y^j \gamma_{ikj}$$

where

$$\gamma_{ikj} = \frac{1}{2} \left(\frac{\partial g_{kj}}{\partial x^i} + \frac{\partial g_{ik}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^k} \right).$$

By asking

$$\gamma_{ij}^{k} = g^{kl} \gamma_{ilj}$$
 i, j, k, l \in {1,...,n},

we have

$$G^k = \frac{1}{2} y^i y^j \gamma^k_{ij} \cdot$$

We note $\Gamma_i^j(x, y) = y^l \Gamma_{il}^j(x)$, the horizontal projector is written:

$$h(\frac{\partial}{\partial x^{i}}) = \frac{\partial}{\partial x^{i}} - \Gamma_{i}^{j} \frac{\partial}{\partial y^{j}}, \ h(\frac{\partial}{\partial y^{j}}) = 0.$$

The vertical projector becomes

$$v(\frac{\partial}{\partial x^i}) = \Gamma_i^j \frac{\partial}{\partial y^j}, \ v(\frac{\partial}{\partial y^i}) = \frac{\partial}{\partial y^i}$$

The curvature $R = \frac{1}{2}[h, h]$ is then

$$R = \frac{1}{2} R_{ij}^k dx^i \wedge dx^j \otimes \frac{\partial}{\partial y^k}$$

with

$$R_{ij}^{k} = \frac{\partial \Gamma_{i}^{k}}{\partial x^{j}} - \frac{\partial \Gamma_{j}^{k}}{\partial x^{i}} + \Gamma_{i}^{l} \frac{\partial \Gamma_{j}^{k}}{\partial y^{l}} - \Gamma_{j}^{l} \frac{\partial \Gamma_{i}^{k}}{\partial y^{l}}, i, j, k, l \in \{1, ..., n\}.$$

Proposition 1

Let E be an energy function, Γ a connection such that $\Gamma\text{=}[J,S].$ The following two relationships are equivalent:

i.
$$i_S dd_J E = -dE$$
;
ii. $d_{\mu}E = 0$.

Proof: We notice that if we have the relation (i), we get

$$i_{s}i_{s}dd_{t}E = 0 = L_{s}E.$$
 (2.6)

We can write successively, taking into account that E is homogeneous of degree 2 and that JS=C,

$$i_{S}dd_{J}E = (L_{S} - di_{S})d_{J}E = L_{S}d_{J}E - 2dE = -dE$$

or
$$L_{S}d_{J}E = dE.$$
 (2.7)

From the formula $L_S d_J - d_J L_S = d_{[S,J]}$, we have

$$L_S d_J E = d_{[S,J]} E + d_J L_S E = (2d_h E - dE) + d_J L_S E.$$

According to the two relations (2.6) and (2.7), we find

$$d_h E = 0.$$

Conversely, if we have $d_h E = 0$, by definition $h = \frac{[J,S] + I}{2}$, we can write

$$d_{[J,S]}E + dE = 0.$$

By noting that $\,hS=S$ and $L_{\rm S}E=0\,$, we can go up the calculation and one obtains

 $i_S dd_J E = -dE$

Properties of the Curvature \mathfrak{R} of a Riemannian Manifold

From the properties of the connection stated above, we obtain the following classic results: for all $X, Y, Z, T \in \chi(TM)$

$$g(\Re(X,Y)Z,JT) = -g(\Re(X,Y)T,JZ);$$
(3.1)

$$\Re(X,Y)Z + \Re(Y,Z)X + \Re(Z,X)Y = 0;$$
(3.2)

$$g(\Re(T,X)Y,JZ) + g(\Re(T,Y)Z,JX) + g(\Re(T,Z)X,JY) = 0;$$
(3.3)

$$g(\Re(X,Y)Z,JT) = g(\Re(Z,T)X,JY).$$
(3.4)

Properties of the Horizontal Nullity Space of Curvatures

From the properties of the curvatures given above and by posing:

$$N_{\mathcal{R}} = \{X \in \chi(TM) \text{ such that } R(X,Y) = 0, \forall Y \in \chi(TM)\},\$$

$$N_{\mathfrak{R}} = \{X \in \chi(TM) \text{ such that } \mathfrak{R}(X,Y)Z = 0, \forall Y, Z \in \chi(TM)\};\$$

we have

$$JX \perp ImR^{\circ} \Leftrightarrow X \in KerR^{\circ} \left(R^{\circ} = i_{S}R \right), X \in \chi(TM);$$

$$IX \perp ImR \Leftrightarrow \mathfrak{M} \left(S \mid X \right) X = 0, X \in \chi(TM) \quad \forall X \in \chi(TM):$$

$$(4.1)$$

$$IX \perp IMK \Leftrightarrow \mathcal{H}(S, X)I = 0, X \in \chi(IM), \forall I \in \chi(IM);$$

$$(4.2)$$

$$X \in N_{R} \Leftrightarrow JX \perp \Re(S, Y)Z, X \in \chi(TM), \forall Y, Z \in \chi(TM);$$
(4.3)

$$X \in N_{\mathfrak{R}} \Leftrightarrow JX \perp \mathfrak{R}(Y, Z)T, X \in \chi(TM), \forall Y, Z, T \in \chi(TM);$$
(4.4)

$$X \in N_{\mathfrak{R}} \Leftrightarrow \mathfrak{R}(Y, Z) X = 0, \forall Y, Z \in \chi(TM).$$

$$(4.5)$$

Properties of the Nullity Space and the Curvature Image Space

Let H=Imh the horizontal space, V=Imv the vertical space and ImR (respectively Im \mathfrak{R}) the module on F(TM) generated by the image of the curvature R (respectively \mathfrak{R}).

Proposition 2

On a Riemannian manifold (M, E) the following properties are verified:

- 1) Im \Re and JN_{\Re} are orthogonal and Im $\Re \oplus JN_{\Re} = V$;
- 2) ImR and $JN_{P}+\{C\}$ are orthogonal; ImR $\oplus \{JN_{P}+\{C\}\}=V$;
- $N_R = N_{\Re}$

4) At the neighborhood of a point z of TM, the subspaces $N_R \cap H, N_{\mathfrak{R}} \cap H, N_R, N_{\mathfrak{R}}, hN_R \oplus JN_R$, ImR and Im \mathfrak{R} are involutive

Proof: The property 1) follows from the relation (4.4).

According to the relation (4.2), JX \perp ImR if and only if $\Re(S, X)Y = 0, \forall Y \in \chi(TM)$. From the relation (2.4), we have

$$\Re(S, X)Y = J[Y, R^{\circ}(X)] - [JY, R^{\circ}(X)] + R([JY, S], X) + R^{\circ}([JY, X]) = 0.$$

Taking into account the relation (4.1), the relation $J[JY,S] = JY, \forall Y \in \chi(TM)$ and that the curvature R is semi-basic, we get

$$R(X,Y) = R^{\circ}[JY,X], \forall Y \in \chi(TM)$$
(5.1)

As R is alternating bilinear on the vector fields of $\chi(TM)$ considered as $\mathcal{F}(TM)$ -module, the relation is not possible unless X=S or, if $X \in N_{R'}$ and if X is generated by projectable vector fields in the nullity space of R. The converse is immediate according to the relations (4.2), (4.3) and (3.3).

The property $N_R = N_{\Re}$ follows from the link between \Re and R given by the relation (2.4) and from the above remark that N_R is generated by projectable vector fields in N_R . At the neighborhood of a point z of TM, the subspaces $N_R \cap H$, $N_{\Re} \cap H$, are involutive. These are well known results see [7]. Let us now show that hN_R \oplus JN_R is involutive. You just have to check on the generators of hN_R and JN_R. Since hN_R is generated by projectable vector fields in hN_R and, hN_R is involutive, then JN_R is generated by the commutative generators in JN_R. JN_R is therefore involutive. It is the same for [X, JY] with $X \in hN_R$ and $Y \in JN_R$. Let K, L \in H such that JK and JL be orthogonal to $JN_{\mathfrak{R}}$. As we have D_{JZ} g=0 for all Z \in H, and taking into account $J[JZ, X] \in JN_{\mathfrak{R}}$ for all $X \in N_{\mathfrak{R}}$, we have

g(J[JK, L], JX)=0, g(J[K, JL], JX)=0,

that is to say, according to the nullity of the Nijenhuis tensor of J, we have $\label{eq:say}$

 $[JK, JL]=J[JK, L]+J[K, JL] \in JN_{ss}$

This proves that $\operatorname{Im} \mathfrak{R}\,$ is involutive.

The space ImR is also involutive, this follows from the assertion 3) and from the following property:

$$g([JY, JZ], C) = 0$$
, for all $JY, JZ \in \text{Im } R$.

Proposition 3

Let ${\sf X}$ be a projectable vector field. The following two relationships are equivalent:

(i) [hX, J]=0

(ii) [JX, h]=0

Proof: For a connection Γ =[J,S], the torsion is zero, that is,

[J,h](X,Y)=0

for all $X, Y \in \chi(TM)$. By developing the above relation, we have,

[JX, hY]+[hX, JY]+Jh[X,Y]+hJ[X,Y]-J[hX, Y]-J[X, hY]-h[JX,Y]-h[X, JY]=0

as we have hJ=0 and, h[X,JY]=0, because X being projectable vector fields by hypothesis, we have

[JX, h]Y+[hX, J]Y=J[X, h]Y

The vector field X being projectable and Jh=J, the term J[X, h] is null. We obtain

[JX, h]Y+[hX, J]Y = 0

Hence the equivalence of the relations (i) and (ii).

Proposition 4

If the rank of the horizontal nullity space hN_R of the curvature R is constant, there is a local base of hN_R verifying the proposition 3.

Proof: We will solve the equation

[hX, J]=0

The vector fields which annihilate the tangent structure by the Lie derivative is well known [8]. The equation to be solved is written in natural local coordinates on an open set $U \subset M$, $(x^i, y^i) \in TU$, $i, j \in \{1, ..., n\}$

$$X^{i}(x)\frac{\partial}{\partial x^{i}} - X^{i}(x)\Gamma_{i}^{j}\frac{\partial}{\partial y^{j}} = X^{i}(x)\frac{\partial}{\partial x^{i}} + y^{i}\frac{\partial X^{j}(x)}{\partial x^{i}}\frac{\partial}{\partial y^{j}}$$

that is to say:

$$y^{i} \frac{\partial X^{j}(x)}{\partial x^{i}} = -X^{i}(x)\Gamma_{i}^{j}$$

As we have $y^{l}\Gamma_{il}^{j}(x) = \Gamma_{i}^{j}(x, y)$, and $\Gamma_{il}^{j} = \Gamma_{li}^{j}$, the previous equation is equivalent to

$$\frac{\partial X^{j}(x)}{\partial x^{i}} = -X^{l}(x)\Gamma_{li}^{j}.$$

The condition for compatibility of the equation, according to the Frobenius theorem, is

$$X^{l}\left(\frac{\partial \Gamma_{li}^{j}}{\partial x^{k}}-\frac{\partial \Gamma_{lk}^{j}}{\partial x^{i}}+\Gamma_{li}^{s}\Gamma_{sk}^{j}-\Gamma_{lk}^{s}\Gamma_{si}^{j}\right)=0, \quad i,j,k,l,s\in\{1,...,n\}.$$

This condition is none other than $X \in hN_{R}$. Hence the result.

Proposition 5

On a Riemannian manifold (M, E), the space H \oplus ImR is involutive if the rank of the space hN $_{\rm R}$ is constant.

Proof: It is clear that for all $X, Y \in \chi(TM)$,

[hX, hY]=h[hX,hY]+v[hX, hY]=h[hX, hY]+R(X,Y)

So $[hX, hY] \in H \oplus ImR.$

According to the proposition 2, it suffices to show that v[hX, R(Y,Z)] is orthogonal to $JN_{R}+\{C\}$. We have

$$D_{hX}g(R(Y,Z),C) - g(D_{hX}R(Y,Z),C) - g(R(Y,Z),D_{hX}C) = 0.$$

The relation (2.3) gives

 $D_{hX}R(Y,Z) = v[hX, R(Y,Z)]$ and $D_{hX}C = [h, C]X$.

The homogeneity of h leads to [h, C]=0, we get

$$g(v[hX, R(Y, Z)], C) = 0.$$

Let $T \in hN_{R}$, then we can write

$$D_{hX}g(R(Y,Z),JT) - g(D_{hX}R(Y,Z),JT) - g(R(Y,Z),D_{hX}JT) = 0.$$

As we have $g(R(Y,Z),JT) = 0, D_{hX}R(Y,Z) = v[hX,R(Y,Z)]$ and $D_{hX}JT = [h, JT]X$, we obtain

$$g(v[hX, R(Y, Z)], JT) + g(R(Y, Z), v[h, JT]X) = 0.$$

According to the proposition 4, we have $v[h, JT]X \in JN_{p}$, therefore

g(v[hX, R(Y, Z)], JT) = 0.

Hence the result.

Proposition 6

On a Riemannian manifold (M, E), the space H \oplus ImR \oplus JN $_{\rm R}$ is involutive if the rank of the space hN $_{\rm p}$ is constant.

Proof: The relation (2.4) is written:

ℜ(X, Y)Z=J[Z, R(X, Y)]-[JZ, R(X, Y)]+R([JZ, X], Y)+R(X, [JZ, Y]).

If we have $Z \in hN_R$, we get $\Re(X, Y)Z=0$ according to the relation (4.5). As the elements of hN_R are generated by projectable vector fields in hN_R , the relation above shows that $[JZ, R(X,Y)] \in ImR \oplus JN_R$, for all $X, Y \in \chi(TM)$. The proposition 5 shows that $H \oplus ImR$ is involutive. We will now show that $v[JZ, hX] \in JN_R$. According to the proposition 4, we can find a base of hN_R verifying the proposition 3. This proves that $v[JZ, hX] \in JN_R$.

Hence the result.

Proposition 7

Let Γ be a linear connection without torsion satisfying the following relationships:

1) there is an energy function E_{a} which satisfies $d_{B}E_{a} = 0$;

2) the space $H \oplus ImR$ is involutive;

3) the vertical space V decomposes into V = ImR \oplus {JN_R +{C}} and the rank of JN_R is constant, then the 1-form $d_v E_o$ is completely integrable

Proof: The kernel of $d_v E_o$ is H the horizontal space of Γ since we have $v \circ h = 0$, ImR by hypothesis, and the vertical vector fields such that $L_{av}E_o = 0$. We will show that these vector fields are involutive.

We can write

[hX, hY]=h[hX, hY]+v[hX, hY]=h[hX, hY]+R(X, Y)

So

 $[hX,hY] \in \!\! H \oplus ImR$

Let v_{\circ} be the vertical projector of the linear connection given by $E_{\circ}, d_{v}E_{\circ}$ and $d_{v_{\circ}}E_{\circ}$ are identical on the vertical space JH since $vJ = v_{\circ}J = J$. The difference between Γ_{\circ} linear connections coming from the energy function E_{\circ} and Γ is in the horizontal space Im h_{\circ} and Imh, the two connections have the same curvature. It remains to calculate v[JY, hX] for all $X, Y \in \chi(TM)$. If JY \in ImR, v[JY, hX] \in ImR by hypothesis. If JY \in JN_R, according to the proposition 4, v[JY, hX] \in JN_R and that H \oplus ImR \oplus JN_R is also involutive according to the proposition 6. The scalar 1-form $d_{v}E_{\circ}$ is completely integrable like the scalar 1-form $d_{v}E_{\circ}$.

Theorem 1

Let Γ be a linear connection without torsion such that the rank of the curvature R is constant. The connection Γ comes from an energy function if and only if

1) the distribution H \oplus ImR is involutive and the vertical space V=ImR \oplus {JN_{_R}+{C}} with hN_{_R} is generated by projectable vector fields in hN_{_R}

2) there is an energy function E_{\circ} such that $d_{R}E_{\circ} = 0$.

Then, there exist a constant ϕ function on the bundles such that $e^{\varphi(x)}E_{\circ}$ is the energy function of Γ

Proof: A necessary and sufficient condition for a connection Γ to come from an energy function E is according to the proposition 1,

d_bE=0

This results in

 $d_h d_h E = 0 = 2d_B E$

The conditions are therefore necessary according to the previous studies. Conversely, let E_{\circ} be an energy function such that $d_{\rm R}E_{\circ}=0$. We will show that with the hypothesis, there exist a constant ϕ function on the bundles such that

 $d_h(e^{\varphi}E_{\circ})=0.$

The equation is equivalent to

 $d\varphi = -\frac{1}{E}d_h E_{\circ}.$

The condition of integrability of such an equation is

$$d(\frac{1}{E_{\circ}}) \wedge d_h E_{\circ} + \frac{1}{E_{\circ}} dd_h E_{\circ} = 0,$$

that is to say

$$dd_h E_\circ = \frac{dE_\circ}{E_\circ} \wedge d_h E_\circ.$$

According to the proposition 7, $d_{\rm v}E_{\rm o}$ is completely integrable. We have, according to the Frobenius theorem,

$$dd_{v}E_{\circ}\wedge d_{v}E_{\circ}=0,$$

By applying the inner product i_c to the above equality, we obtain

$$dd_{v}E_{\circ} = \frac{dE_{\circ}}{E_{\circ}} \wedge d_{v}E_{\circ},$$

that is to say

$$dd_h E_\circ = \frac{dE_\circ}{E_\circ} \wedge d_h E_\circ.$$

This is the condition of integrability sought.

Let M be a differentiable, paracompact, connected manifold of dimension n \geq 2, Γ a linear connection without torsion such that the rank of the curvature is locally constant over M. Then, Γ is a canonical connection of a Riemannian manifold if and only if at the neighborhood of any point where the rank of the curvature R is constant;

1) The distribution $H \oplus ImR$ is involutive and V=ImR $\oplus \{JN_R + \{C\}\}\$ with hN_n is generated by projectable vector fields in hN_n

2) There is an energy function E_{a} such that $d_{B}E_{a} = 0$

Proof: It follows from the theorem 1. Using a partition of the unit, we glue together the local metrics to have a global metric.

The search for an energy function E of Γ leads to the search for an energy function E_{\circ} such that $d_{R}E_{\circ} = 0$. In natural local coordinates, the curvature R is written:

$$R = \frac{1}{2} y^{l} R_{l,ij}^{k}(x) dx^{l} \wedge dx^{j} \otimes \frac{\partial}{\partial y^{k}}, i, j, k, l \in \{1, ..., n\}.$$

An energy function is written

$$E_{\circ}=\frac{1}{2}g_{ij}^{\circ}y^{i}y^{j}.$$

Thus, the relation $d_{\rm \scriptscriptstyle R} E_{\rm \scriptscriptstyle o} = 0$ is equivalent to the following system of equations:

$$g_{kl}^{\circ}R_{l,ij}^{k} = 0$$

$$g_{kl}^{\circ}R_{r,ij}^{k} = -g_{kr}^{\circ}R_{l,ij}^{k} \text{ with } l \neq r.$$

In the matrix form, the system is written:

$$\begin{bmatrix} g_{11}^{\circ} \cdots g_{n1}^{\circ} \\ \vdots & \ddots & \vdots \\ g_{1n}^{\circ} \cdots g_{nn}^{\circ} \end{bmatrix} \begin{pmatrix} R_{1,j}^{1} \cdots R_{n,jj}^{1} \\ \vdots & \ddots & \vdots \\ R_{1,j}^{n} \cdots R_{n,jj}^{n} \end{bmatrix} = \begin{bmatrix} 0 \\ \ddots & -^{t} A \\ A \\ \ddots & 0 \end{bmatrix}$$

Remark: Solving the inverse problem leads to the system of linear equations, the other conditions are easy to verify, which is quite easy, in practice. In addition, the trace of $R_{k,ii}^k$ is zero.

On an Isotropic Finslerian Manifold [9]

An isotropic finslerian manifold is a manifold [10] with a connection such that

$$R^{\circ} = 2EKJ - Kd_{I}E \otimes C,$$

 $R^{\circ} = i_{\circ}R_{\gamma}K$ denotes the sectional curvature.

If $K \neq 0$, the manifold is of regular curvature, that is to say, the dimension of ImR is equal to n-1. Our theorem remains valid. In the case of a finslerian manifold of dimension 2, the theorem is written.

Theorem 2 [11]

Let M be a connected paracompact differentiable manifold of dimension 2, Γ a connection without torsion with non-zero curvature over a dense subset of TM=TM-{0}.

 Γ comes from a finslerian structure if and only if, at the neighborhood of any point where the curvature is non-zero,

i. H \oplus ImR is involutive

ii. The contracted of the curvature \tilde{R} is such $d\tilde{R}$ is of maximum rank and the function $i_{s}\tilde{R}$ positive

To illustrate the results, we will give some examples.

Example 1

Let the manifolds $M=\mathbb{R}^3$ have coordinates $(x^1, x^2, x^3) \in \mathbb{R}^3$, (y^1, y^2, y^3) on the bundle of the tangent space $T\mathbb{R}^3$, Γ the connection [J, S] with

$$S = y^{1} \frac{\partial}{\partial x^{1}} + y^{2} \frac{\partial}{\partial x^{2}} + y^{3} \frac{\partial}{\partial x^{3}} - 2((y^{1})^{2} e^{x^{3}} + y^{2} y^{3}) \frac{\partial}{\partial y^{1}}.$$

The coefficients Γ_i^j of Γ being $\Gamma_i^j = \frac{\partial G^j}{\partial y^i}$, the non-zero coefficients are

$$\Gamma_1^1 = 2y^1 e^{x^3}, \Gamma_2^1 = y^3, \Gamma_3^1 = y^2.$$

A base of the horizontal space of Γ is written

$$\frac{\partial}{\partial x^{1}} - 2y^{1}e^{x^{3}}\frac{\partial}{\partial y^{1}}, \frac{\partial}{\partial x^{2}} - y^{3}\frac{\partial}{\partial y^{1}}, \frac{\partial}{\partial x^{3}} - y^{2}\frac{\partial}{\partial y^{1}}$$

The curvature R is written:

$$R = e^{x^3} dx^1 \wedge ((y^1 - y^2) dx^3 - y^3 dx^2) \otimes \frac{\partial}{\partial y^1}.$$

The horizontal nullity space is generated by

$$(y^{1}-y^{2})\frac{\partial}{\partial x^{2}}+y^{2}\frac{\partial}{\partial x^{3}}-[(y^{2})^{2}+(y^{1}-y^{2})y^{3}]\frac{\partial}{\partial y^{1}}$$

The horizontal nullity space is not generated by projectable vector fields in $hN_{\rm R}$, $N_{\rm R}$ is not involutive. This linear connection according to the proposition 2 cannot come from an energy function.

Example 2

We take $M = \mathbb{R}^4$ and the energy function is written:

$$E = \frac{1}{2} (e^{x^{1}} (y^{1})^{2} + e^{x^{2}} (y^{2})^{2} + e^{x^{3}} (y^{3})^{2} + e^{2x^{1}} (y^{4})^{2}).$$

The canonical spray of E is written:

$$S = y^{1} \frac{\partial}{\partial x^{1}} + y^{2} \frac{\partial}{\partial x^{2}} + y^{3} \frac{\partial}{\partial x^{3}} + y^{4} \frac{\partial}{\partial x^{4}} - (\frac{1}{2}(y^{1})^{2} - (y^{4})^{2}e^{x^{1}})\frac{\partial}{\partial y^{1}} - \frac{1}{2}(y^{2})^{2} \frac{\partial}{\partial y^{2}} - \frac{1}{2}(y^{3})^{2} \frac{\partial}{\partial y^{3}}$$

The near zero coefficients of E are

The non-zero coefficients of Γ are

$$\Gamma_1^1 = \frac{1}{2}y^1, \Gamma_4^1 = -y^4 e^{x^1}, \Gamma_2^2 = \frac{1}{2}y^2, \Gamma_3^3 = \frac{1}{2}y^3, \Gamma_1^4 = y^4, \Gamma_4^4 = y^1.$$

The horizontal fields are generated by

$$\frac{\partial}{\partial x^{i}} - \frac{1}{2}y^{i}\frac{\partial}{\partial y^{1}} - y^{4}\frac{\partial}{\partial y^{4}}, \frac{\partial}{\partial x^{2}} - \frac{1}{2}y^{2}\frac{\partial}{\partial y^{2}}, \frac{\partial}{\partial x^{3}} - \frac{1}{2}y^{3}\frac{\partial}{\partial y^{3}}, \frac{\partial}{\partial x^{4}} + y^{4}e^{x^{i}}\frac{\partial}{\partial y^{1}} - y^{1}\frac{\partial}{\partial y^{4}}$$

The horizontal nullity space is generated by

$$\frac{\partial}{\partial x^2} - \frac{1}{2} y^2 \frac{\partial}{\partial y^2}, \frac{\partial}{\partial x^3} - \frac{1}{2} y^3 \frac{\partial}{\partial y^3}$$

The two basic elements of the nullity space verifying the proposition 3 are:

$$e^{\frac{-x^2}{2}}(\frac{\partial}{\partial x^2}-\frac{1}{2}y^2\frac{\partial}{\partial y^2}), e^{\frac{-x^3}{2}}(\frac{\partial}{\partial x^3}-\frac{1}{2}y^3\frac{\partial}{\partial y^3}).$$

The non-zero elements of the curvature are:

$$R_{4,14}^{1} = -\frac{e^{x^{1}}}{2}, R_{1,14}^{4} = \frac{1}{2}.$$

and Im $R = \left\langle -y^{4}e^{x^{1}}\frac{\partial}{\partial y^{1}} + y^{1}\frac{\partial}{\partial y^{4}} \right\rangle \mathcal{F}(\mathsf{T}\mathbb{R}^{4})$

We see that ${\rm H} \oplus {\rm ImR}$ is involutive.

The matrix is written

e^{x^1}	0	0	0	0	0	0	$-\frac{e^{x^{l}}}{2}$		0	0	0	$-\frac{e^{2x^{l}}}{2}$
0	e^{x^2}	0	0	0	0	0	0		0	0	0	0
0	0	e^{x}	0	0	0	0	0	-	0	0	0	0
0	0	0	e^{2x^1}	$\left\ \frac{1}{2}\right\ $	0	0	0		$\left \frac{e^{2x^{l}}}{2}\right $	0	0	0

Example 3

We take $M=\mathbb{R}^4$ and the energy function is written:

$$E = \frac{1}{2} (e^{x^4} (y^1)^2 + e^{x^4} (y^2)^2 + e^{x^3} (y^3)^2 + (y^4)^2).$$

The canonical spray of E is written:

$$S = y^{1}\frac{\partial}{\partial x^{1}} + y^{2}\frac{\partial}{\partial x^{2}} + y^{3}\frac{\partial}{\partial x^{3}} + y^{4}\frac{\partial}{\partial x^{4}} - y^{1}y^{4}\frac{\partial}{\partial y^{1}} - y^{2}y^{4}\frac{\partial}{\partial y^{2}} - \frac{1}{2}(y^{3})^{2}\frac{\partial}{\partial y^{3}} + \frac{e^{x^{4}}((y^{1})^{2} + (y^{2})^{2})}{2}\frac{\partial}{\partial y^{4}}$$

The non-zero coefficients of Γ are

$$\Gamma_{1}^{1} = \frac{1}{2}y^{4}, \Gamma_{4}^{1} = \frac{1}{2}y^{1}, \Gamma_{2}^{2} = \frac{1}{2}y^{4}, \Gamma_{4}^{2} = \frac{1}{2}y^{2}, \Gamma_{3}^{3} = \frac{1}{2}y^{3}, \Gamma_{1}^{4} = -\frac{e^{x^{4}}y^{1}}{2}, \Gamma_{2}^{4} = -\frac{e^{x^{4}}y^{2}}{2}$$

The horizontal fields are generated by

$$\frac{\partial}{\partial x^{1}} - \frac{1}{2}y^{4}\frac{\partial}{\partial y^{1}} + \frac{e^{x^{4}}y^{1}}{2}\frac{\partial}{\partial y^{4}}, \frac{\partial}{\partial x^{2}} - \frac{1}{2}y^{4}\frac{\partial}{\partial y^{2}} + \frac{e^{x^{4}}y^{2}}{2}\frac{\partial}{\partial y^{4}}, \frac{\partial}{\partial x^{3}} - \frac{1}{2}y^{3}\frac{\partial}{\partial y^{3}},$$
$$\frac{\partial}{\partial x^{4}} - \frac{1}{2}y^{1}\frac{\partial}{\partial y^{1}} - \frac{1}{2}y^{2}\frac{\partial}{\partial y^{2}}.$$

The horizontal nullity space is generated by

$$\frac{\partial}{\partial x^3} - \frac{1}{2} y^3 \frac{\partial}{\partial y^3}.$$

The basic element of the nullity space verifying the proposition 3 is:

$$e^{\frac{-x^3}{2}}(\frac{\partial}{\partial x^3}-\frac{1}{2}y^3\frac{\partial}{\partial y^3}).$$

The non-zero elements of the curvature are:

$$R_{2,12}^{1} = -\frac{e^{x^{4}}}{4}, R_{4,14}^{1} = -\frac{1}{4}, R_{1,12}^{2} = \frac{e^{x^{4}}}{4},$$
$$R_{4,24}^{2} = -\frac{1}{4}, R_{1,14}^{4} = \frac{e^{x^{4}}}{4}, R_{2,24}^{4} = \frac{e^{x^{4}}}{4}$$

And ImR generated by

$$\left\langle e^{x^4} \left(-y^2 \frac{\partial}{\partial y^1} + y^1 \frac{\partial}{\partial y^2} \right), -y^4 \frac{\partial}{\partial y^1} + y^1 e^{x^4} \frac{\partial}{\partial y^4}, -y^4 \frac{\partial}{\partial y^2} + y^2 e^{x^4} \frac{\partial}{\partial y^4} \right\rangle$$

 $\mathcal{F}(\mathbb{T}\mathbb{R}^4).$

We see that $H \oplus ImR$ is involutive.

The matrix is written

Conclusion

To recognize a connection without torsion, that is to say,

 $\Gamma = [J, S] and [C, S] = S$

comes from an energy function.

Case 1

The connection is of regular curvature, that is to say, the module generated on $\mathcal{F}(\text{TM})$ by ImR is of dimension n-1. The necessary and sufficient condition is written:

- i. $H \oplus ImR$ is involutive
- ii. There is an energy function E_{\circ} such that $d_R E_{\circ} = 0$

Indeed, according to proposition 1, we have $d_h E=0$, this implies $d_R E=0$. So $H \oplus ImR$ is the kernel of dE. Then, $H \oplus ImR$ is involutive, and $d_v E_\circ$ is a completely integrable 1-form.

This result is valid for a Finslerian manifold.

Case 2

dim(ImR)<n-1. The necessary and sufficient condition becomes

i. The horizontal nullity space of the curvature R is generated by projectable vector fields in $hN_{_R}$, the space H \oplus ImR is involutive and the vertical space is written V=ImR \oplus {JN}_{_R}+{C}}

ii. There exists an energy function E_{\circ} such that $d_{R}E_{\circ} = 0$

Indeed, it is essential to have the horizontal nullity space of the curvature generated by projectable vector fields in hN_R . This assertion involves that N_R is involutive. With the help of the metric g, we have $H \oplus ImR$ is involutive, the scalar 1-form $d_v E_s$ is then completely integrable.

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