

Geometric Formula for Prime form on a Sewn Riemann Surface

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Abstract

We use the geometric representation for the Szegő kernel on genus g_1+g_2 Riemann surface obtained in the frames of the Yamada sewing construction of two Riemann surfaces of genus g_1 and genus g_2 in order to derive new formulas expressing prime forms. These formulas can be used in vertex operator algebra and conformal field theory computation as well as in algebraic geometry.

Keywords: Szegő kernel; Prime forms; Theta-functions

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Differentials and Kernels on a Riemann Surface

Differentials on a Riemann surface

Consider a compact Riemann surface $\Sigma^{(g)}$ of genus g with canonical homology cycle basis $a_1, \dots, a_g, b_1, \dots, b_g$. There exists g holomorphic one-forms $v_i^{(g)}, i=1, \dots, g$ which may be normalized [1,2] by:

$$\oint_{a_i} v_j^{(g)} = 2\pi i \delta_{ij}. \quad (1)$$

The genus g period matrix $\Omega^{(g)}$ is defined by:

$$\Omega_{ij}^{(g)} = \frac{1}{2\pi i} \oint_{b_j} v_i^{(g)}, \quad (2)$$

For $i, j=1, \dots, g$, $\Omega^{(g)}$ is symmetric with positive imaginary part, i.e., $\Omega^{(g)} \in \mathbb{H}_g$, the Siegel upper half plane. Next we give the definition of the theta function with real characteristics [1,3].

$$\theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] (z | \Omega^{(g)}) = \sum_{m \in \mathbb{Z}^g} \exp(i\pi(m + \alpha^{(g)})\Omega^{(g)}(m + \alpha^{(g)}) + (z + 2\pi i\beta^{(g)})),$$

for $\alpha^{(g)} = (\alpha_i), \beta^{(g)} = (\beta_i) \in \mathbb{R}^g, z = (z_i) \in \mathbb{C}^g$, and $i=1, \dots, g$. There exists [3,4] a non-singular odd character $\left[\begin{matrix} \gamma^{(g)} \\ \delta^{(g)} \end{matrix} \right]$ such that

$$\theta^{(g)} \left[\begin{matrix} \gamma^{(g)} \\ \delta^{(g)} \end{matrix} \right] (0 | \Omega^{(g)}) = 0, \partial_{z_i} \theta^{(g)} \left[\begin{matrix} \gamma^{(g)} \\ \delta^{(g)} \end{matrix} \right] (0 | \Omega^{(g)}) \neq 0,$$

Then introduce:

$$\zeta^{(g)}(x) = \sum_{i=1}^g \partial_{z_i} \theta^{(g)} \left[\begin{matrix} \gamma^{(g)} \\ \delta^{(g)} \end{matrix} \right] (0 | \Omega^{(g)}) v_i^{(g)}(x), \quad (3)$$

a holomorphic one-form, and let $(\zeta^{(g)}(x))^{\frac{1}{2}}$ denote the form of weight $\frac{1}{2}$ on the double cover $\widetilde{\Sigma^{(g)}}$ of $\Sigma^{(g)}$. We also refer to $(\zeta^{(g)}(x))^{\frac{1}{2}}$ as a double-valued $\frac{1}{2}$ -form on $\Sigma^{(g)}$.

Let us define the prime form $\mathcal{E}^{(g)}(x, y)$ by:

$$\mathcal{E}^{(g)}(x, y) = \theta^{(g)} \left[\begin{matrix} \gamma^{(g)} \\ \delta^{(g)} \end{matrix} \right] \left(\int_y^x v^{(g)} | \Omega^{(g)} \right) (\zeta^{(g)}(x))^{-\frac{1}{2}} (\zeta^{(g)}(y))^{\frac{1}{2}}, \quad (4)$$

where $\int_y^x v^{(g)} = \left(\int_y^x v_i^{(g)} \right) \in \mathbb{C}^g$. The prime form $\mathcal{E}^{(g)}(x, y) = -\mathcal{E}^{(g)}(y, x)$ is a

holomorphic differential form of weight $(-\frac{1}{2}, -\frac{1}{2})$ on $\widetilde{\Sigma^{(g)}} \times \widetilde{\Sigma^{(g)}}$.

We define the Szegő kernel [4-6] for $\theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] (0 | \Omega^{(g)}) \neq 0$ by the formula:

$$S^{(g)} \left[\begin{matrix} \theta^{(g)} \\ \phi^{(g)} \end{matrix} \right] (x, y | \Omega^{(g)}) = \frac{\theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] \left(\int_y^x v^{(g)} | \Omega^{(g)} \right)}{\theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] (0 | \Omega^{(g)}) \mathcal{E}^{(g)}(x, y)}, \quad (5)$$

where $\theta^{(g)} = (\theta_i), \phi^{(g)} = (\phi_i) \in U(1)^n$ for $\theta_j = -e^{-2\pi i \beta_j}, \phi_j = -e^{-2\pi i \alpha_j}, j=1, \dots, g$. This can be written as:

$$S^{(g)} \left[\begin{matrix} \theta^{(g)} \\ \phi^{(g)} \end{matrix} \right] (x, y | \Omega^{(g)}) = \Theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] (x, y; v^{(g)}, 0 | \Omega^{(g)}) (\mathcal{E}^{(g)}(x, y))^{-1},$$

with the functional:

$$\Theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] (x, y; f_1, f_2 | \Omega^{(g)}) = \frac{\theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] \left(\int_y^x f_1 | \Omega^{(g)} \right)}{\theta^{(g)} \left[\begin{matrix} \alpha^{(g)} \\ \beta^{(g)} \end{matrix} \right] (f_2 | \Omega^{(g)})}.$$

The Szegő Kernel in the ε -formalism

In Appendix we recall the ε -formalism due to Yamada [7] of sewing of two Riemann surfaces of genera g_1 and g_2 to form a genus g_1+g_2 Riemann surface. In a study [8] we determined the Szegő kernel on the Riemann surface $\Sigma^{(g_1+g_2)}$ in terms of data coming from Szegő kernel:

$$S^{(g_a)}(x, y) = S^{(g_a)}[\theta^{(g_a)}] \quad (6)$$

on the surface $\Sigma^{(g_a)}$ for $a=1,2$. Let us recall that construction here. We adopt the abbreviated notation of the left hand side of eqn. (6)

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when there is no ambiguity. Similarly, the Szegő kernel on $\Sigma^{(g_1+g_2)}$ is denoted by:

$$S^{(g_1+g_2)}(x, y) = S^{(g_1+g_2)}[\theta^{(g_1+g_2)}] \quad (7)$$

with periodicities $(\theta_{s_a}^{(g_1+g_2)}, \theta_{s_a}^{(g_1+g_2)}) = (\theta_{s_a}^{(g_a)}, \theta_{s_a}^{(g_a)})$ on the inherited homology basis. Note that we exclude those Riemann theta characteristics for which eqn. (7) exists but where either of the lower genus theta functions vanishes, i.e., we assume that eqn. (6) exists for $a=1, 2$.

Following a study [8] we define weighted moments for $S^{(g_1+g_2)}$ by:

$$X_{ab}(k, l, \varepsilon) = X_{ab}[\theta^{(g_1+g_2)}] = \frac{\varepsilon^{\frac{1}{2}(k+l-1)}}{(2\pi i)^2} \oint_{c_a(x)c_b(y)} x^{-k} y^{-l} S^{(g_1+g_2)}(x, y) dx^{\frac{1}{2}} dy^{\frac{1}{2}}, \quad (8)$$

for $k, l \geq 1$. It follows that:

$$X_{ab}[\theta^{(g_1+g_2)}] \quad (9)$$

We denote by $X_{ab} = (X_{ab}(k, l, \varepsilon))$ the infinite matrix indexed by $k, l \geq 1$.

We also define various moments for $S^{(g_a)}(x, y)$. These provide the data used to construct $S^{(g_1+g_2)}(x, y)$. Define holomorphic $\frac{1}{2}$ -forms on $\hat{\Sigma}^{(g_a)}$ by:

$$h_a(k, x, \varepsilon) = h_a[\theta^{(g_a)}] \quad (10)$$

$$\bar{h}_a(k, y, \varepsilon) = \bar{h}_a[\theta^{(g_a)}] \quad (11)$$

and introduce infinite row vectors $h_a(x) = (h_a(k, x))$, $h_a(y) = (h_a(k, y))$, indexed by $k \geq 1$. It follows that:

$$\bar{h}_a[\theta^{(g_a)}] \quad (12)$$

Finally, we define the moment matrix:

$$\begin{aligned} F_a(k, l, \varepsilon) &= F_a[\theta^{(g_a)}] = \frac{\varepsilon^{\frac{1}{2}(k+l-1)}}{(2\pi i)^2} \oint_{c_a(x)c_a(y)} x^{-k} y^{-l} S^{(g_a)}(x, y) dx^{\frac{1}{2}} dy^{\frac{1}{2}} \\ &= \frac{\varepsilon^{\frac{k}{2} - \frac{1}{4}}}{2\pi i} \oint_{c_a(x)} x^{-k} h_a(l, x) dx^{\frac{1}{2}} = \frac{\varepsilon^{\frac{l}{2} - \frac{1}{4}}}{2\pi i} \oint_{c_a(y)} y^{-l} \bar{h}_a(k, y) dy^{\frac{1}{2}}. \end{aligned} \quad (13)$$

$F_a(k, l, \varepsilon)$ obeys a skew-symmetry property similar to eqn. (9).

We are now in a position to express $S^{(g_1+g_2)}(x, y)$ in terms of the lower genus data. From the sewing relation eqn. (23) we have

$$dz_a = -\varepsilon \frac{dz_{\bar{a}}}{z_{\bar{a}}^2} \text{ so that:}$$

$$dz_a^{\frac{1}{2}} = (-1)^{\bar{a}} \xi \varepsilon^{\frac{1}{2}} \frac{dz_{\bar{a}}^{\frac{1}{2}}}{z_{\bar{a}}}, \quad (14)$$

where $\xi \in \{\pm\sqrt{-1}\}$ determines the square root branch chosen. We then find:

Proposition 1: $S^{(g_1+g_2)}(x, y)$ is given by:

$$S^{(g_1+g_2)}(x, y) = \begin{cases} S^{(g_a)}(x, y) + h_a(x) X_{\bar{a}\bar{a}} \bar{h}_a^T(y), & \text{for } x, y \in \hat{\Sigma}^{(g_a)}, \\ h_a(x) (\xi(-1)^{\bar{a}} I - X_{\bar{a}\bar{a}}) \bar{h}_a^T(y), & \text{for } x \in \hat{\Sigma}^{(g_a)}, y \in \hat{\Sigma}^{(g_{\bar{a}})}, \end{cases} \quad (15)$$

Where I denotes the infinite identity matrix and T denotes the transpose.

In a study [8] we computed the explicit form of the moment matrix X_{ab} in terms of the moments F_a of $S^{(g_a)}(x, y)$. It is useful to introduce infinite block matrices:

$$X = \begin{pmatrix} X_{11} & X_{12} \\ X_{21} & X_{22} \end{pmatrix}, \quad F = \begin{pmatrix} F_1 & 0 \\ 0 & F_2 \end{pmatrix}, \quad \Xi = \begin{pmatrix} 0 & \xi I \\ -\xi I & 0 \end{pmatrix}, \quad Q = F\Xi = \begin{pmatrix} 0 & \xi F_1 \\ -\xi F_2 & 0 \end{pmatrix}. \quad (16)$$

Then one finds:

Proposition 2: X is given by:

$$X = (I - Q)^{-1} F \quad (17)$$

where $(I - Q)^{-1} = \sum_{n \geq 0} Q^n$ is convergent for $|\varepsilon| < |r_1 r_2|$.

Propositions 1 and 2 imply:

Theorem 1: $S^{(g_1+g_2)}(x, y)$ is given by:

$$S^{(g_1+g_2)}(x, y) = \delta_{ab} S^{(g_a)}(x, y) + h_a(x) (\Xi(I - Q)^{-1})_{ab} \bar{h}_b^T(y),$$

For $x \in \hat{\Sigma}^{(g_a)}, y \in \hat{\Sigma}^{(g_b)}$. Equivalently,

$$S^{(g_1+g_2)}(x, y) = \begin{cases} S^{(g_a)}(x, y) + h_a(x) (I - F_a F_a)^{-1} F_a \bar{h}_a^T(y), & \text{for } x, y \in \hat{\Sigma}^{(g_a)}, \\ \xi(-1)^{\bar{a}} h_a(x) (I - F_a F_a)^{-1} \bar{h}_a^T(y), & \text{for } x \in \hat{\Sigma}^{(g_a)}, y \in \hat{\Sigma}^{(g_{\bar{a}})}. \end{cases} \quad (18)$$

Similarly to a study [9] we define the determinant of $I - Q$ as a formal power series in $\varepsilon^{\frac{1}{2}}$ by:

$$\log \det(I - Q) = \text{Tr} \log(I - Q) = - \sum_{n \geq 1} \frac{1}{n} \text{Tr}(Q^n).$$

Clearly $\text{Tr}(Q^{2k}) = 2 \text{Tr}((F_1 F_2)^k)$ for $k \geq 0$ whereas $\text{Tr}(Q^n) = 0$ or n odd. Furthermore, from eqn. (13) the diagonal terms $(F_1 F_2)^k$ have integral power series in ε . Thus it follows that [8]:

Lemma 1: $\det(I - Q) = \det(I - F_1 F_2)$ and is a formal power series in ε .

The determinant has the following holomorphic properties:

Theorem 2: $\det(I - Q)$ is non-vanishing and holomorphic in ε or $|\varepsilon| < |r_1 r_2|$.

Geometric Formulas for Prime Forms

In this section we derive formulas for the prime form on a Riemann surface of genus g_1+g_2 sewn of two Riemann surfaces of genera g_1 and g_2 . Let us denote:

$$J_{\bar{a}\bar{a}} = \begin{cases} h_a(x) (\xi(-1)^{\bar{a}} I - X_{\bar{a}\bar{a}}) \bar{h}_a^T(y), & \text{for } x \in \hat{\Sigma}^{(g_a)}, y \in \hat{\Sigma}^{(g_{\bar{a}})}, \\ h_a(x) X_{\bar{a}\bar{a}} \bar{h}_a^T(y), & \text{for } x, y \in \hat{\Sigma}^{(g_a)}. \end{cases}$$

Then we obtain:

Proposition 3

$$\begin{aligned} \mathcal{E}^{(g_1+g_2)}(x, y) &= J_{\bar{a}\bar{a}}^{-1} \left[\Theta^{(g_1+g_2)} \begin{bmatrix} \alpha^{(g_1+g_2)} \\ \beta^{(g_1+g_2)} \end{bmatrix} \right] \left(x, y; \nu^{(g_1+g_2)}, 0 \mid \Omega^{(g_1+g_2)} \right) \\ &\quad - \delta_{\bar{a}\bar{a}} \Theta^{(g_a)} \left[\alpha^{(g_a)} \right] \left(x, y; \nu^{(g_a)}, 0 \mid \Omega^{(g_a)} \right) e^{\frac{1}{2} h_a X_{\bar{a}\bar{a}} h_a^T}, \end{aligned} \quad (19)$$

or, equivalently,

$$\mathcal{E}^{(g_1+g_2)}(x, y) = \Theta^{(g_1+g_2)} \left[\alpha^{(g_1+g_2)} \right] \left(x, y; \nu^{(g_1+g_2)}, 0 \mid \Omega^{(g_1+g_2)} \right)$$

$$\cdot \left[\delta_{a\bar{a}} \left(\begin{bmatrix} \alpha^{(g_a)} \\ \beta^{(g_a)} \end{bmatrix} \right) \left(0 | \Omega^{(g)} \right) \right]^{-1} \left(\zeta^{(g_a)}(x) \right)^{\frac{1}{2}} \left(\zeta^{(g_a)}(y) \right)^{\frac{1}{2}} + J_{a\bar{a}} \right]^{-1}, \quad (20)$$

Proof. In a study [9] it was proved that:

$$\mathcal{E}^{(g_1+g_2)}(x, y) = \mathcal{E}^{(g_a)}(x, y) e^{-\frac{1}{2} b_a X_{a\bar{a}} b_a^T}.$$

Using eqn. (22) we obtain:

$$\Theta^{(g_1+g_2)} \left[\begin{bmatrix} \alpha^{(g_1+g_2)} \\ \beta^{(g_1+g_2)} \end{bmatrix} \right] \left(x, y; \nu^{(g_1+g_2)}, 0 | \Omega^{(g_1+g_2)} \right) \left(\mathcal{E}^{(g_1+g_2)}(x, y) \right)^{-1} \quad (21)$$

$$= \delta_{a\bar{a}} \Theta^{(g_a)} \left[\begin{bmatrix} \alpha^{(g_a)} \\ \beta^{(g_a)} \end{bmatrix} \right] \left(x, y; \nu^{(g_a)}, 0 | \Omega^{(g_a)} \right) \left(\mathcal{E}^{(g_1+g_2)}(x, y) \right)^{-1} e^{-\frac{1}{2} b_a X_{a\bar{a}} b_a^T} + J_{a\bar{a}},$$

and therefore:

$$\mathcal{E}^{(g_1+g_2)}(x, y) = \Theta^{(g_1+g_2)} \left[\begin{bmatrix} \alpha^{(g_1+g_2)} \\ \beta^{(g_1+g_2)} \end{bmatrix} \right] \left(x, y; \nu^{(g_1+g_2)}, 0 | \Omega^{(g_1+g_2)} \right)$$

$$\left[\delta_{a\bar{a}} \Theta^{(g_a)} \left[\begin{bmatrix} \alpha^{(g_a)} \\ \beta^{(g_a)} \end{bmatrix} \right] \left(x, y; \nu^{(g_a)}, 0 | \Omega^{(g_a)} \right) \left(\mathcal{E}^{(g_1+g_2)}(x, y) \right)^{-1} e^{-\frac{1}{2} b_a X_{a\bar{a}} b_a^T} + J_{a\bar{a}} \right]^{-1}.$$

Thus we get eqns. (19) and (20) using and the definition eqn. (4).

Remark 1: Note that in eqns. (19) and (20), $J_{a\bar{a}}$ can be expressed as in eqn. (18), i.e.,

$$J_{a\bar{a}} = \begin{cases} h_a(x) (I - F_a F_a)^{-1} F_a \bar{h}_a^T(y), & \text{for } x, y \in \hat{\Sigma}^{(g_a)}, \\ \xi(-1)^{\bar{a}} h_a(x) (I - F_a F_a)^{-1} \bar{h}_a^T(y), & \text{for } x \in \hat{\Sigma}^{(g_a)}, y \in \hat{\Sigma}^{(g_{\bar{a}})}, \end{cases} \quad (22)$$

where I denotes the infinite identity matrix and T the transpose.

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